

Radiative and Auger Recombination in Indium Nitride from First-Principles



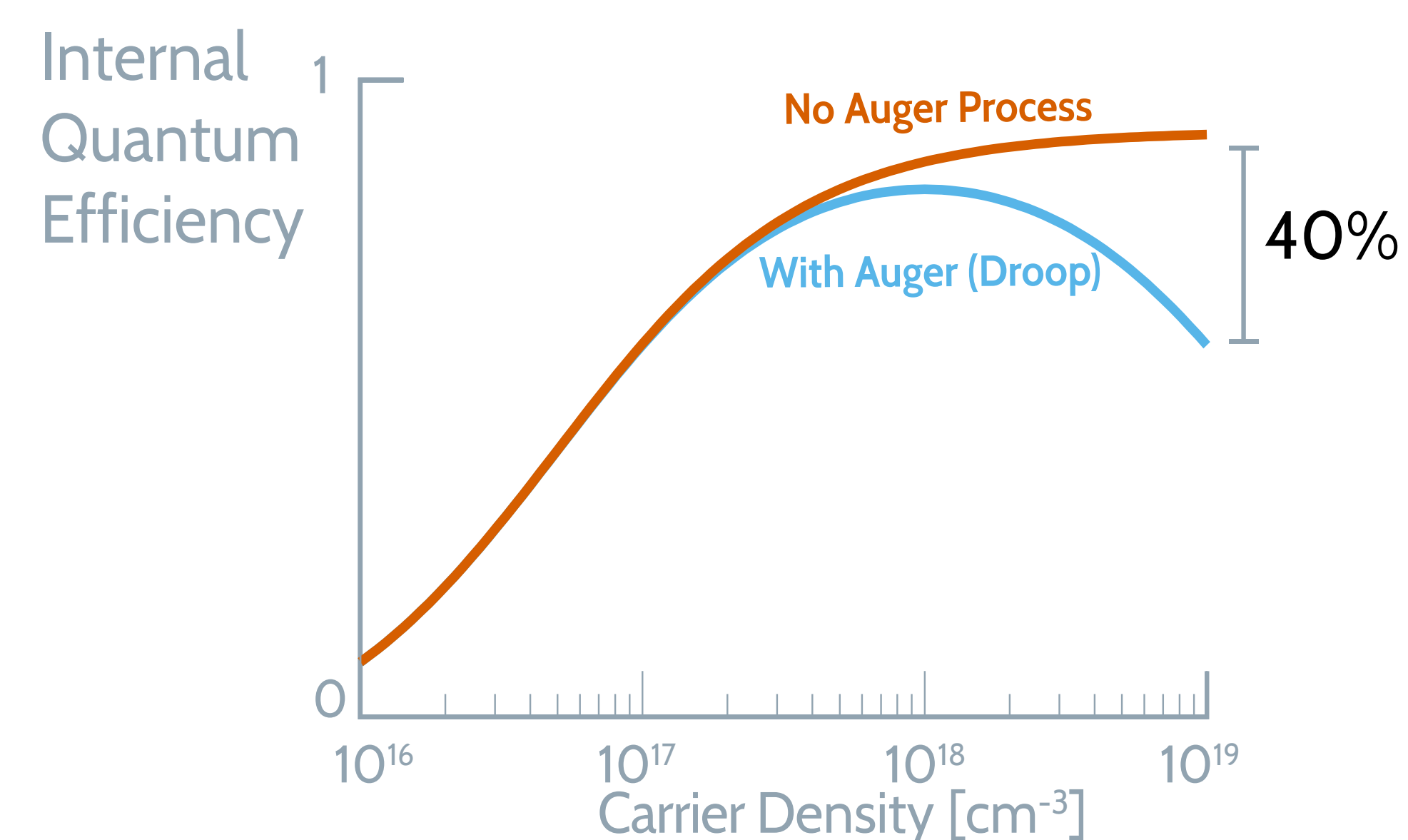
Andrew McAllister¹, Dylan Bayerl², Emmanouil Kioupakis²

¹Applied Physics Program ²Materials Science and Engineering, University of Michigan

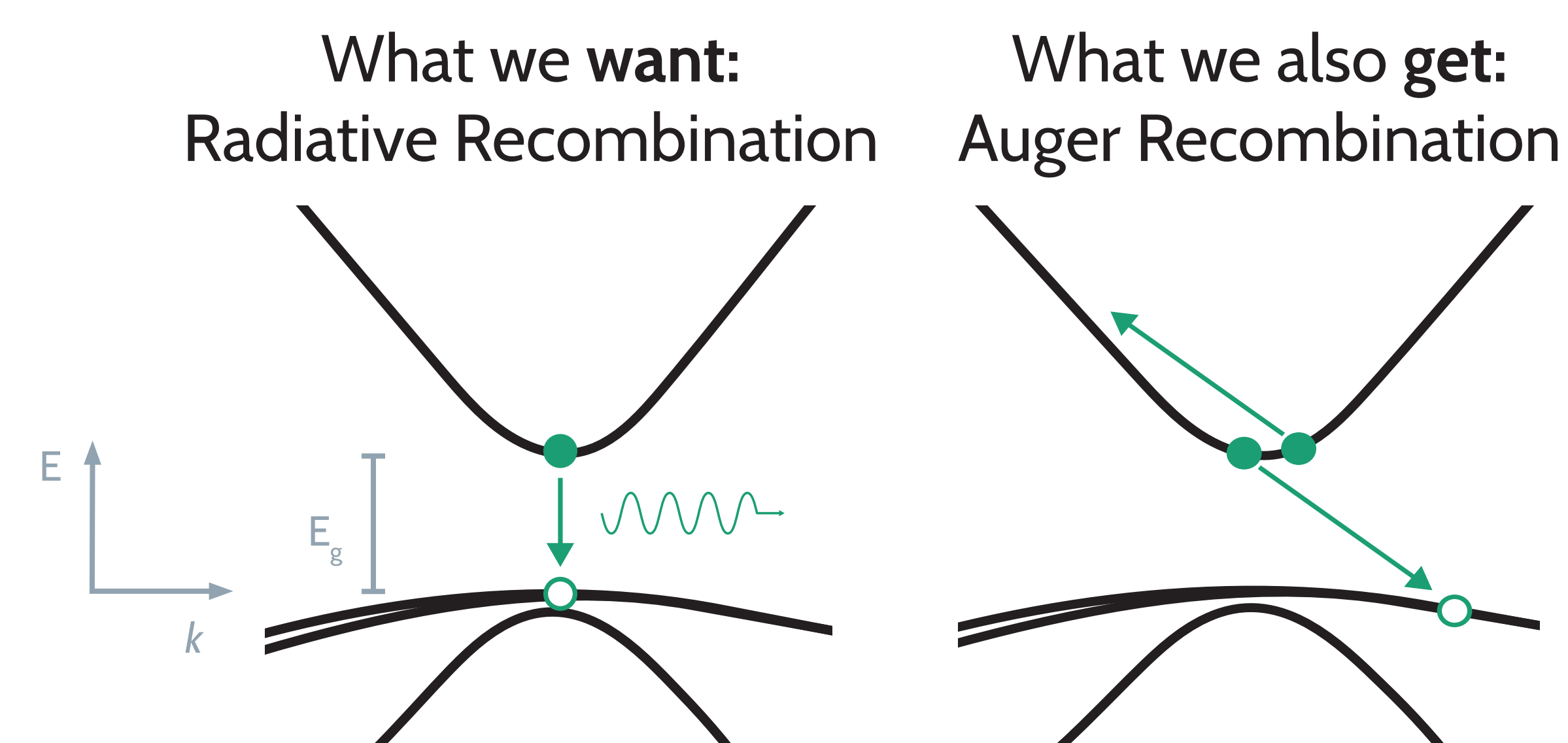
Why Indium Nitride?

The group-III nitrides (AlN, GaN, and InN and alloys) have revolutionized the lighting industry. With gaps ranging from the deep-UV to the infrared, these devices can be tuned to many different emission wavelengths. The center of excitement has been the middle of this range: GaN and Ga-rich InGaN. These materials have allowed the creation of efficient blue light-emitting diodes (LEDs), which efficiently create white light through phosphor down-conversion. But the infrared end of the alloy spectrum could help another industry. InN has a very small gap (0.7 eV) which is near the primary telecommunications wavelength of 1550 nm (0.8 eV). In-rich InGaN could be tuned to this wave length, and the material has other properties for telecom applications (Winden, JJAP, 2013).

The Problem: Efficiency Drop-off or “Droop”

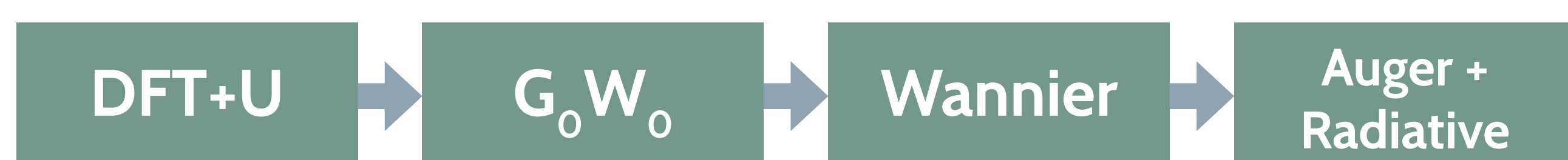


As carrier density increases, efficiency goes down

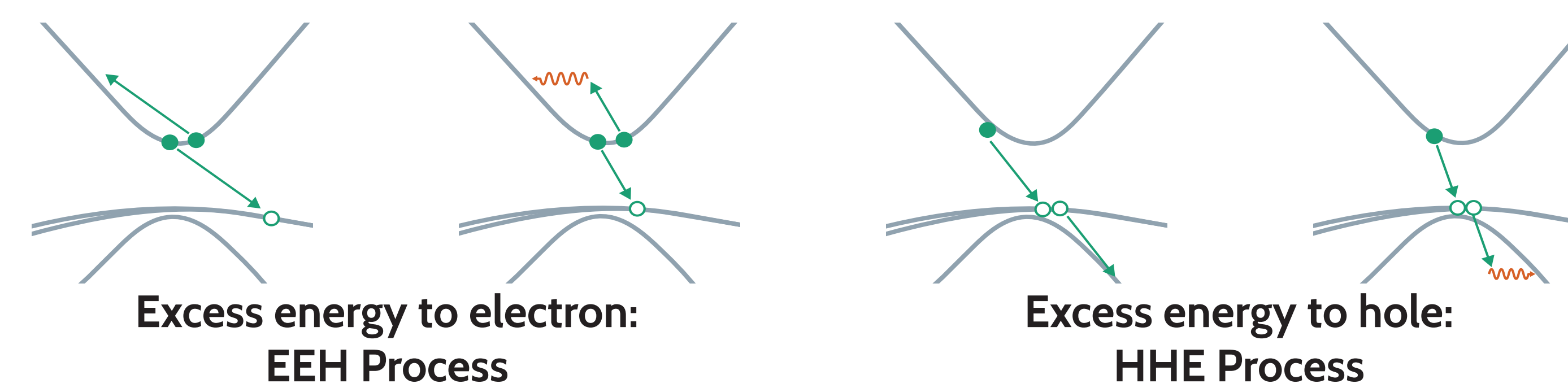


Auger is expected to *increase* as the band gap of a material decreases, making InN more likely to suffer from the Auger process. Because of the non-radiative nature of the process, it is difficult to observe experimentally. The magnitude of its influence on InN is not known with certainty; experimental estimates place the C coefficient between 10^{-28} to 10^{-30} cm⁶s⁻¹ (Cho, APL, 2013; Jang, APL, 2008; Tsai, APL 2007).

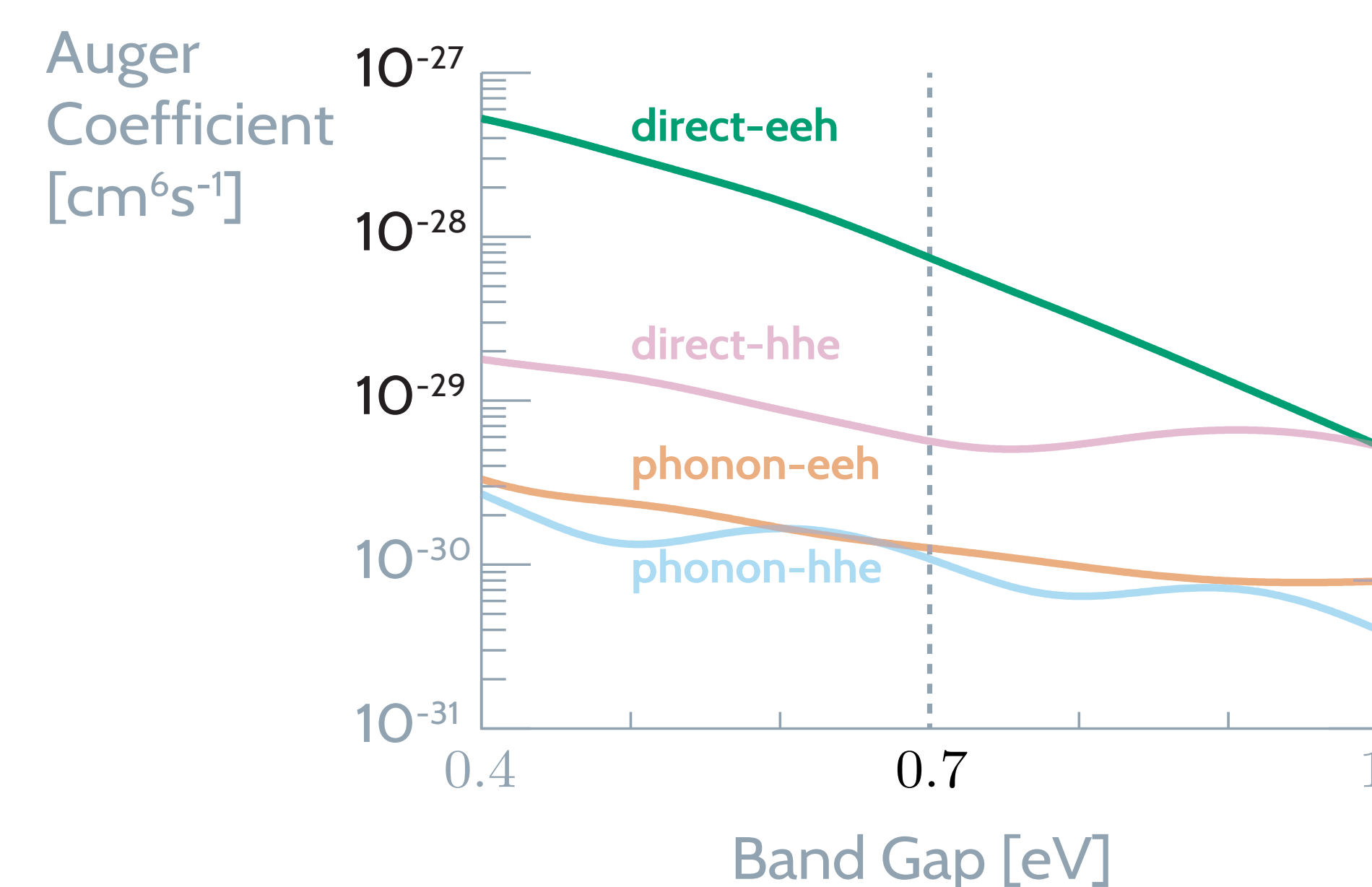
Our First-Principles Workflow



Which Auger Process is Most Important?



Auger recombination comes in different forms. For large gap materials like GaN, assisted processes are dominant (Kioupakis, APL, 2011). However, for a small-gap material we expect to find that *direct Auger* is dominant.



Direct EEH Auger is dominant in indium nitride

Recombination with degenerate carriers

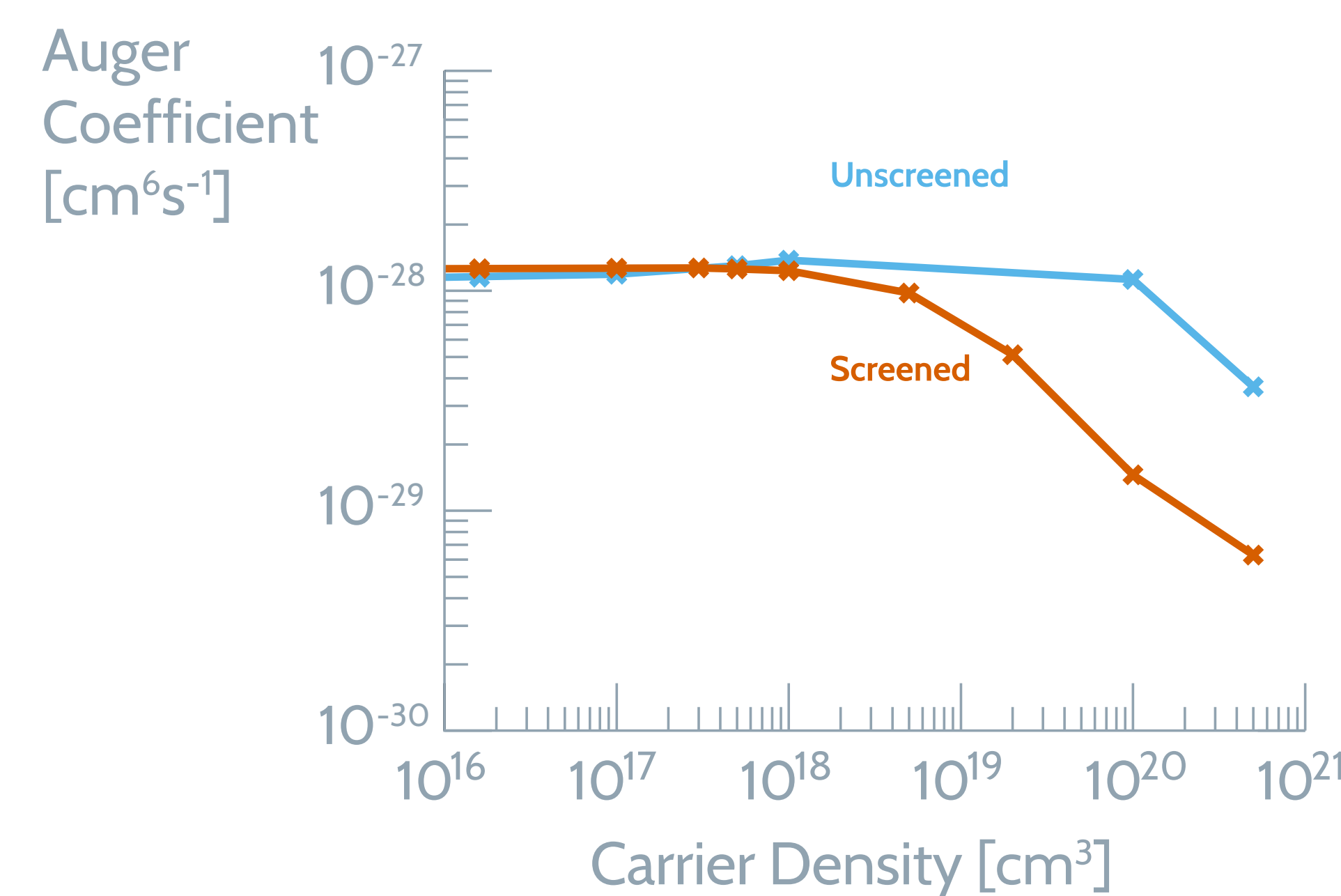
Recombination in semiconductors is typically modeled using the *ABC Model*. At higher carrier densities the model must be modified.

$$\frac{dn}{dt} = R = \underbrace{An}_{\text{Defect}} + \underbrace{Bn^2}_{\text{Radiative}} + \underbrace{Cn^3}_{\text{Auger}} \quad \rightarrow \quad R' = An + B'n + C'n^2$$

Non-degenerate Degenerate

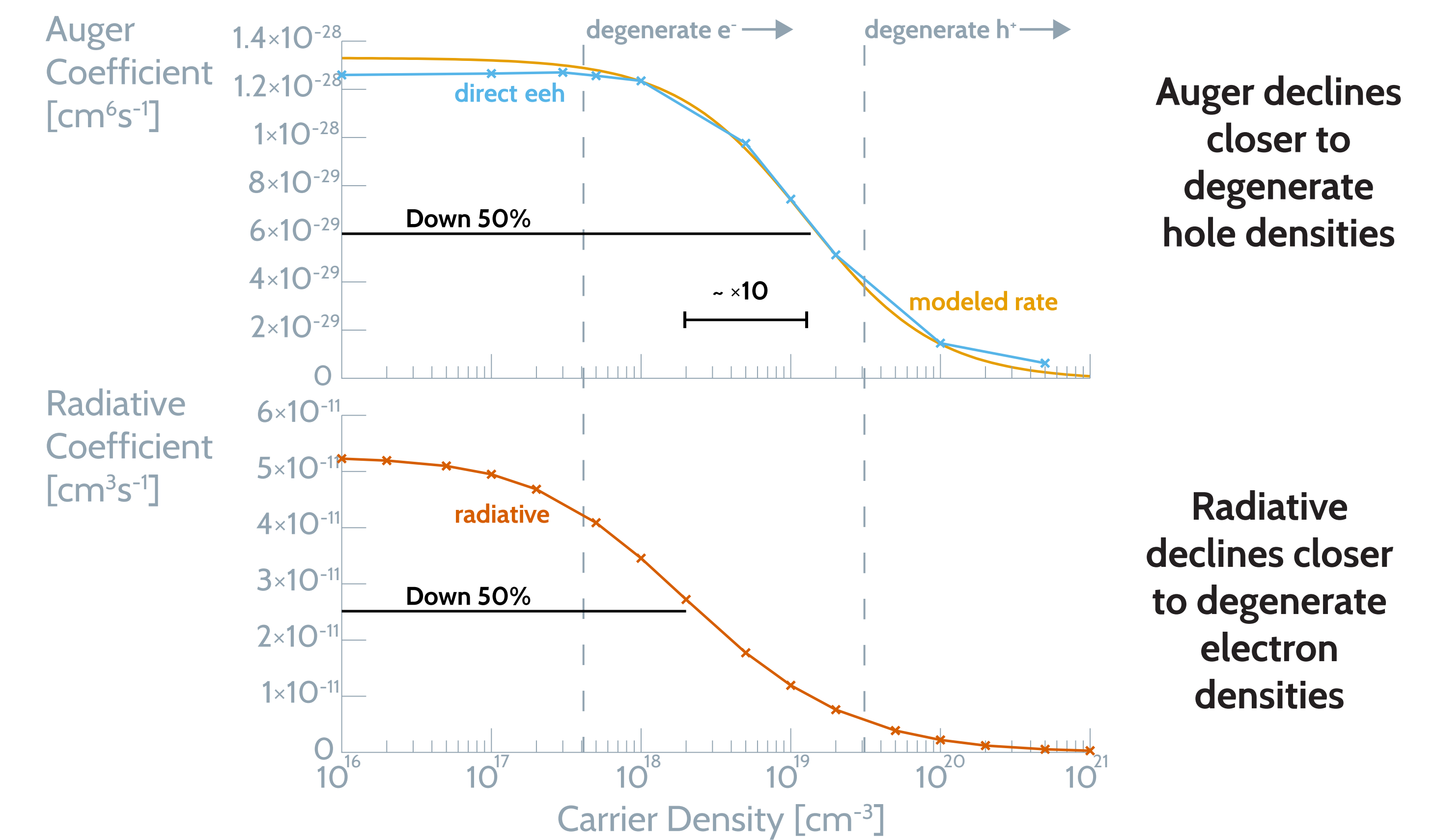
Phase Space Filling: Transitions forbidden by occupied states
Carrier Screening: Carriers not feeling effects of all carriers

Which affects Auger More?



Screening affects Auger rate before Phase-Space Filling

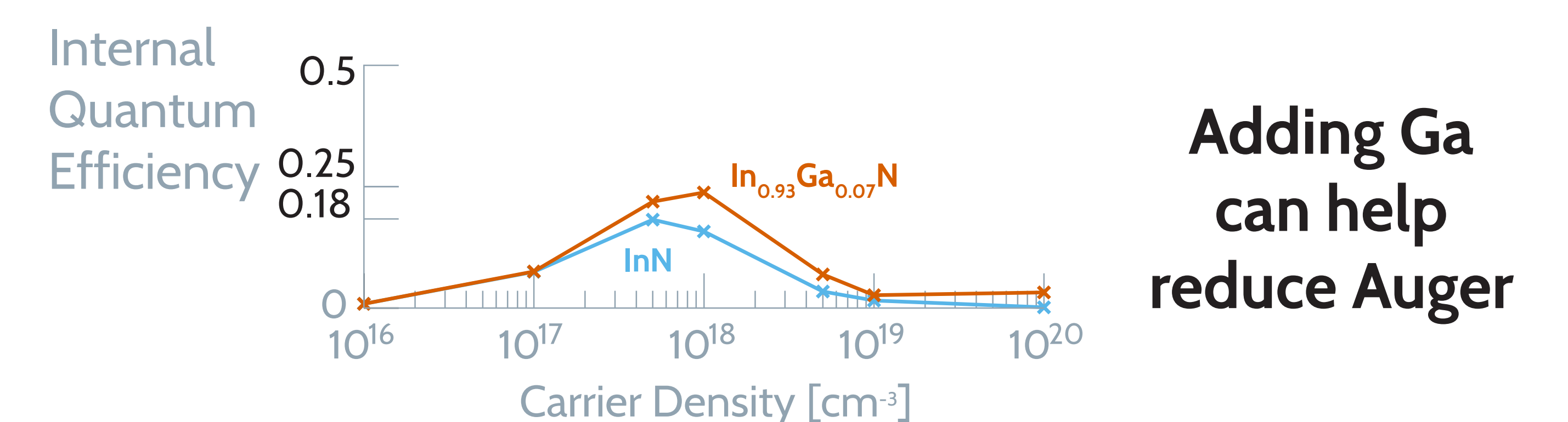
Degeneracy affects the radiative and Auger rates differently



We also modeled the Auger rate as a function of carrier density using:

$$f(x) = \frac{C_0}{1 + \left(\frac{n}{n_0}\right)^b} \quad \begin{array}{l} \text{Low Density Limit} \\ \text{Transition and} \\ \text{Density Dependence} \end{array} \quad \begin{array}{l} \text{At high densities,} \\ \text{Auger is } n^2 \end{array}$$

Improvement of efficiency by alloying with GaN



Adding Ga can help reduce Auger

Because the Auger and radiative processes begin to decline at different carrier densities it becomes a major hurdle to create efficient pure InN devices. But by only changing the gap slightly or reducing the background carrier density, we can create more efficient In-rich InGaN devices with gaps at 0.8 eV. These devices would emit in the main telecommunications range (~1550 nm).

Support From

Funding provided by:
NSF GRFP (G.N. DGE 1256260)
NSF CAREER (G.N. DMR-1254314)

Computational resources provided by:
University of Michigan ARC-TS
NERSC

Developers of:
Quantum Espresso, Wannier 90, Berkeley GW

